



BOUT++ Simulations of Edge Turbulence in the Alcator C-Mod Tokamak

E. M. Davis¹, M. Porkolab¹, J.W. Hughes¹, and X. Q. Xu²

¹Plasma Science and Fusion Center, MIT, Cambridge, MA 02139

²Lawrence Livermore National Laboratory, Livermore, CA 94550

BOUT++ Workshop, Livermore, CA, September 16, 2011

*Work supported by USDoE awards DE-FC02-99-ER54512, DE-FG02-94-ER54235, and DE-AC52-07NA27344, and NNSA SSGF



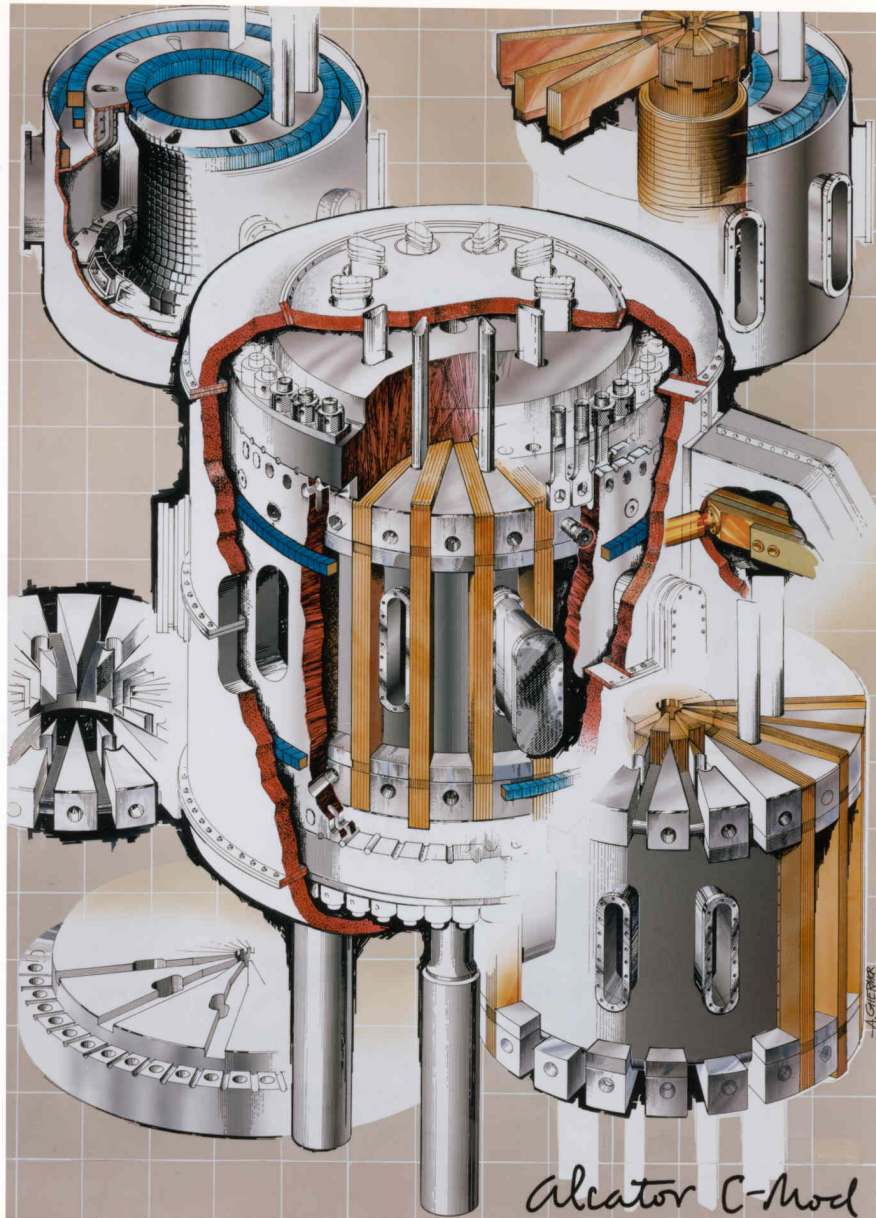
Outline

- Introduction
- Alcator C-Mod Overview
- Physics of the Peeling-Ballooning Module
Sequentially Added to BOUT++
Simulations of C-Mod
 - Resistivity
 - Diamagnetic and $E \times B$ flow
 - Nonlinear Effects
- Conclusion and Future Work

Introduction

- Tokamak energy confinement is thought to be strongly controlled by plasma transport in the edge region just inside the last closed magnetic flux surface
 - A first principles understanding of this transport requires coupling between experiment and theory
- BOUT++ is capable of nonlinear fluid boundary turbulence analysis in a general geometry
 - Experimentally measured C-Mod profiles have been used as input for BOUT++ simulations

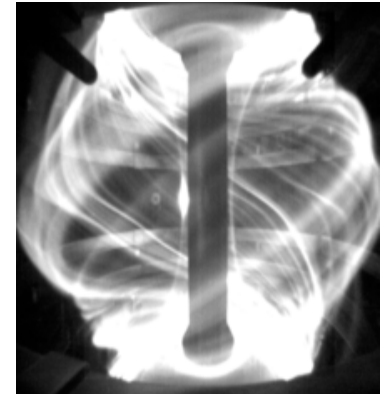
The Alcator C-Mod Tokamak



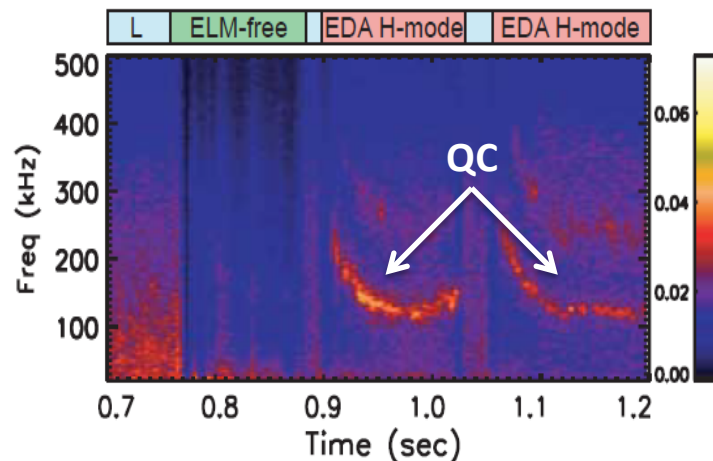
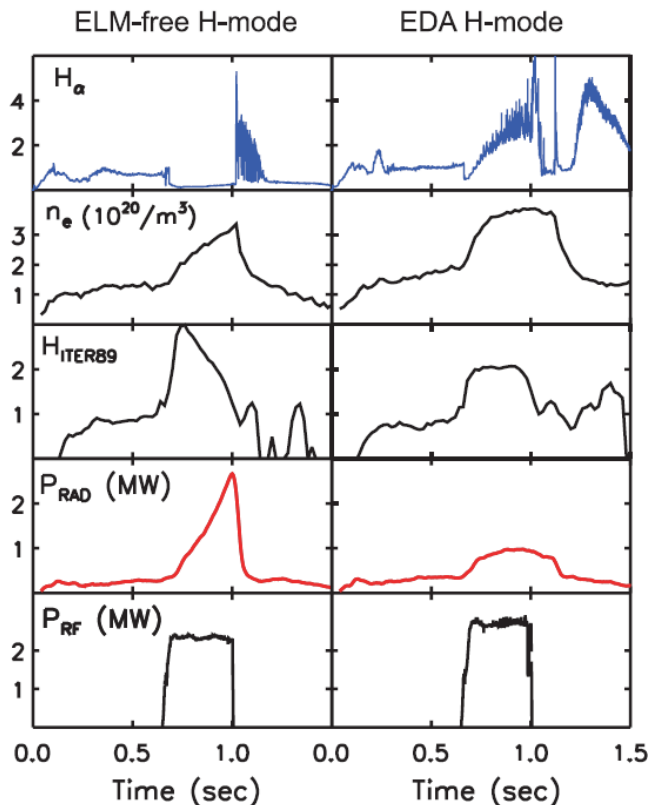
- Alcator C-Mod is a compact, high-field tokamak
 - $R = 0.66 \text{ m}$, $a = 0.22 \text{ m}$
 - $B < 8.1 \text{ T}$
 - $n_e < 1 \times 10^{21} \text{ m}^{-3}$
 - $I_p < 2 \text{ MA}$
- Active Research
 - Heating and Current Drive
 - Plasma Transport
 - Edge and Divertor Physics
- Several confinement regimes are investigated on C-mod
 - H-mode
 - **I-mode**
 - L-mode
 - Linear Ohmic

Edge Transport Strongly Influences Energy Confinement

- L-mode confinement is not likely to lead to a viable fusion reactor
- H-mode confinement is satisfactory for an economic fusion reactor
- Edge Localized Modes (ELMs) or other mild edge modes (Quasi-Coherent or Weakly Coherent Modes) reduce impurity accumulation and allow steady state H mode operation
- I mode is presently under investigation at C-Mod and elsewhere



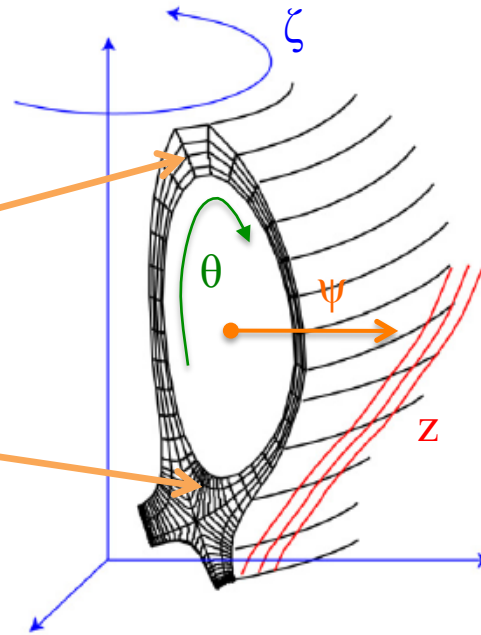
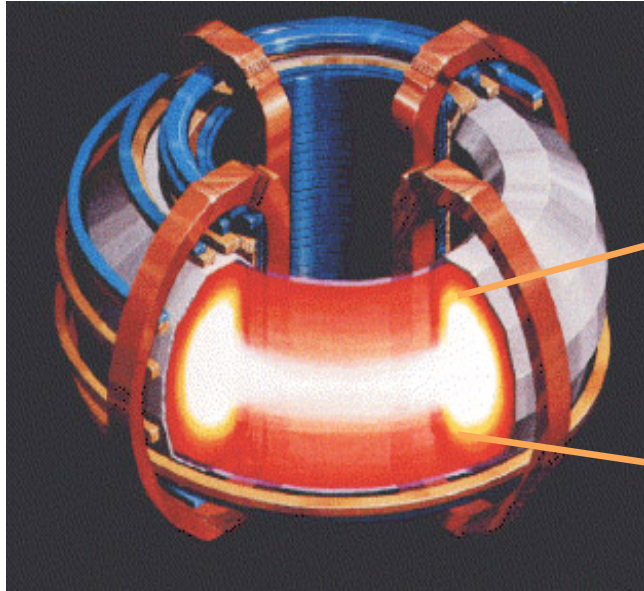
An ELM in MAST



PCI measured density fluctuations in various confinement regimes
* M. Greenwald, et. al, *Fusion Sc. and Tech.*, **51**, 266, (2007)

- C-Mod's **enhanced D_α (EDA) H-mode** is relatively quiescent with good energy confinement and reduced impurity confinement
 - Pedestal regulated by a continuous quasi-coherent mode (**QCM**) oscillation between 50 - 200 kHz

Magnetic geometry in BOUT++ Edge Plasmas



Field-aligned coordinates

$$x = \psi - \psi_0,$$

$$y = \theta,$$

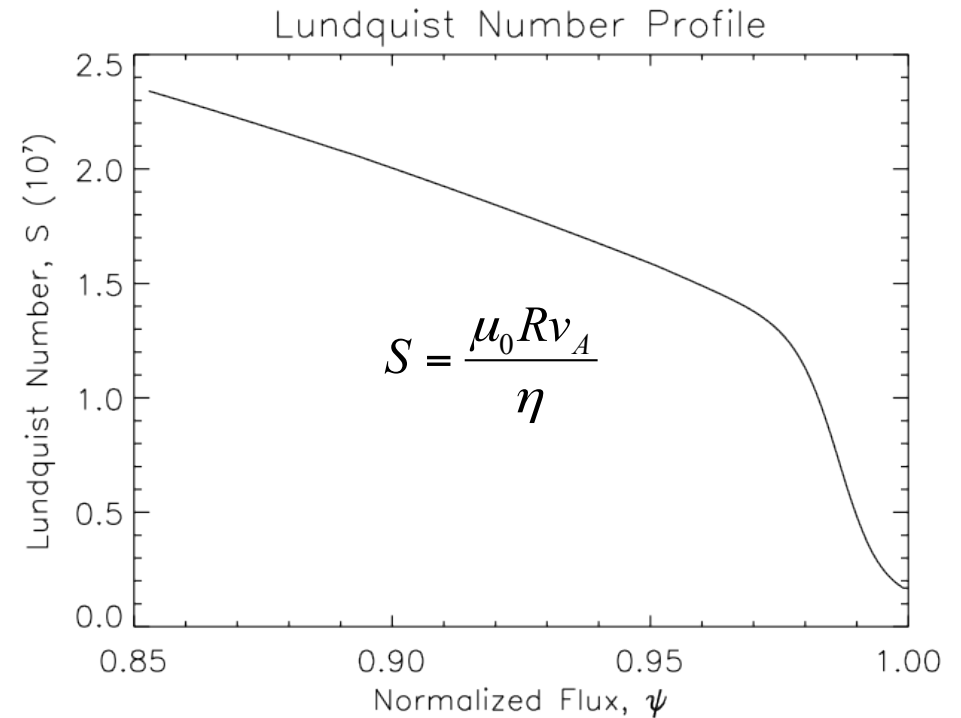
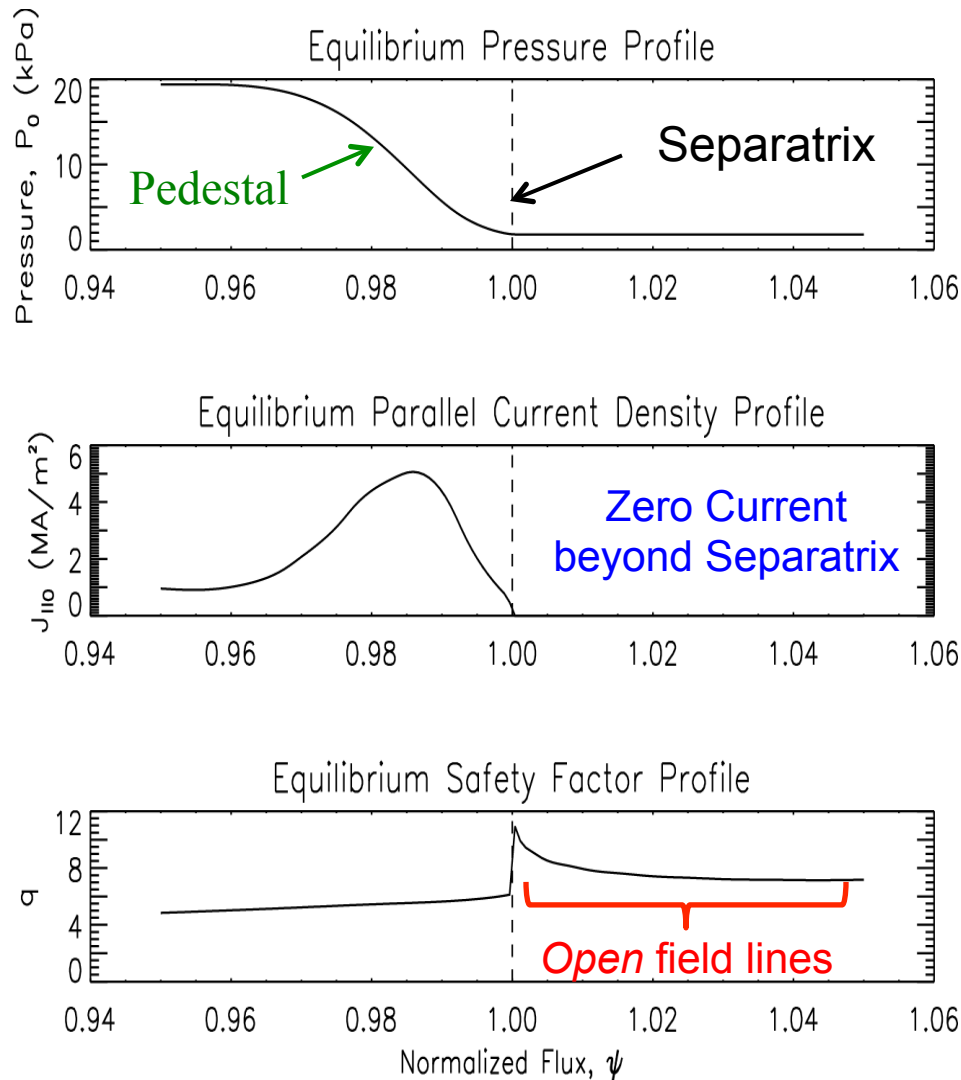
$$z = \zeta - \int_{\theta_0}^{\theta} v(\psi, \theta) d\theta$$

where v is the local safety factor given by:

$$v(\psi, \theta) = \frac{\vec{B} \cdot \nabla \zeta}{\vec{B} \cdot \nabla \theta}$$

- Magnetic field topology changes from closed to open field lines across the separatrix
- In this talk, the edge refers to $0.95 < \psi < 1.05$ (~ 1 cm region in C-Mod)

C-Mod Equilibrium EDA H-Mode Parameters used as BOUT++ Input (1110201023.00900)



Lundquist Number (S) is a dimensionless ratio of the resistive diffusion time to the Alfvén time

– $S \sim 10^7$ in C-Mod EDA pedestal

The Nonlinear System of Equations for Simulating Non-Ideal MHD Peeling-Ballooning Modes

Reduced MHD Equations	Vorticity	$\frac{\partial \varpi}{\partial t} + v_E \cdot \nabla \varpi = B_0^2 \nabla_{\parallel} \left(\frac{j_{\parallel}}{B_0} \right) + 2b_0 \times \kappa \cdot \nabla p,$	Non-ideal physics ✓ Using resistive MHD term, resistivity can be renormalized as Lundquist Number $S = \mu_0 R v_A / \eta$ ✓ Using hyper-resistivity η_H $S_H = \mu_0 R^3 v_A / \eta_H = S / \alpha_H$ ✓ After gyroviscous cancellation, the diamagnetic drift modifies the vorticity and additional nonlinear terms ✓ Using force balance and assuming no net rotation, $E_{r0} = (1/N_i Z_i e) \nabla_{\perp} P_{i0}$
	Pressure	$\frac{\partial P}{\partial t} + v_E \cdot \nabla P = 0,$	
	Ohm's	$\frac{\partial A_{\parallel}}{\partial t} = -\nabla_{\parallel} (\phi + \Phi_0) + \frac{\eta}{\mu_0} \nabla_{\perp}^2 A_{\parallel} - \frac{\eta_H}{\mu_0} \nabla_{\perp}^4 A_{\parallel},$	
Definitions	$\varpi = \frac{n_0 M_i}{B_0} \left(\nabla_{\perp}^2 \phi + \frac{1}{n_0 Z_i e} \nabla_{\perp}^2 p_i \right), \quad P = P_0 + p$		
	$j_{\parallel} = J_{\parallel 0} - \frac{1}{\mu_0} \nabla_{\perp}^2 A_{\parallel}, \quad v_E = \frac{1}{B_0} b_0 \times \nabla (\phi + \Phi_0)$		

Non-Ideal Physics were Methodically Included in Simulations after Initial Ideal Simulations

Reduced MHD Equations	Vorticity	$\frac{\partial \varpi}{\partial t} + v_E \cdot \nabla \varpi = B_0^2 \nabla_{\parallel} \left(\frac{j_{\parallel}}{B_0} \right) + 2b_0 \times \kappa \cdot \nabla p,$
	Pressure	$\frac{\partial P}{\partial t} + v_E \cdot \nabla P = 0,$
	Ohm's	$\frac{\partial A_{\parallel}}{\partial t} = -\nabla_{\parallel} (\phi + \cancel{\Phi_0}) + \frac{\eta}{\mu_0} \nabla_{\perp}^2 A_{\parallel} - \cancel{\frac{\eta_H}{\mu_0} \nabla_{\perp}^4 A_{\parallel}},$

Definitions

$$\varpi = \frac{n_0 M_i}{B_0} \left(\nabla_{\perp}^2 \phi + \cancel{\frac{1}{n_0 Z_i e} \nabla_{\perp}^2 p_i} \right), \quad P = P_0 + p$$

$$j_{\parallel} = J_{\parallel 0} - \frac{1}{\mu_0} \nabla_{\perp}^2 A_{\parallel}, \quad v_E = \frac{1}{B_0} b_0 \times \nabla (\phi + \cancel{\Phi_0})$$

Non-ideal physics

✓ Using resistive MHD term, resistivity can be renormalized as **Lundquist Number**

$$S = \mu_0 R v_A / \eta$$

✓ Using hyper-resistivity η_H

$$S_H = \mu_0 R^3 v_A / \eta_H = S / \alpha_H$$

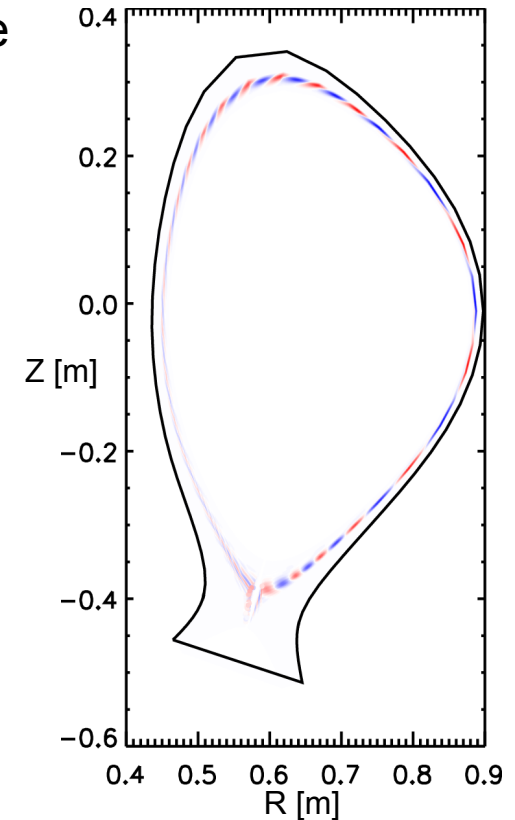
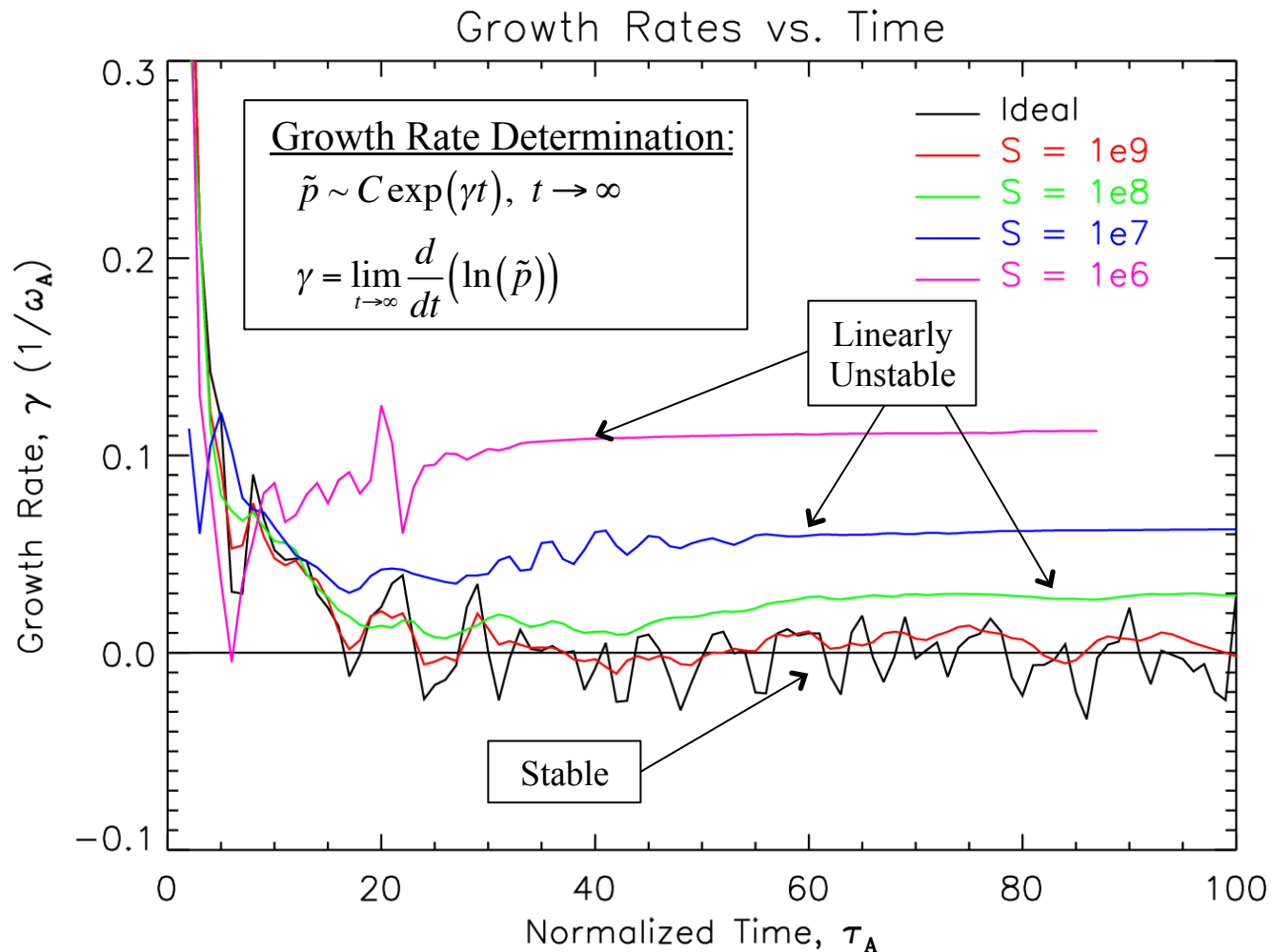
✓ After gyroviscous cancellation, the diamagnetic drift modifies the vorticity and additional nonlinear terms

✓ Using force balance and assuming no net rotation,

$$E_{r0} = (1/N_i Z_i e) \nabla_{\perp} P_{i0}$$

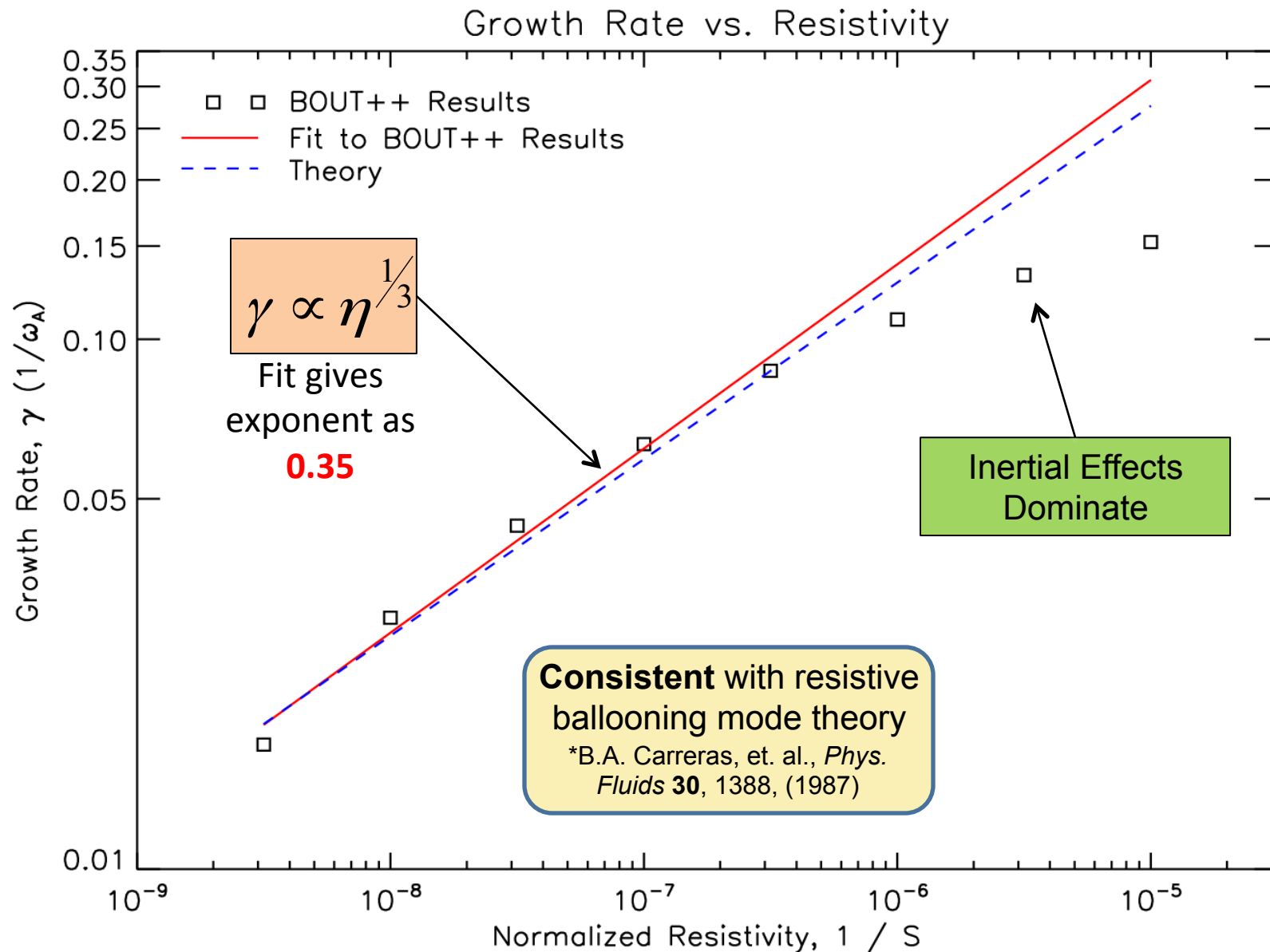
BOUT++ Calculations Show C-Mod EDA H-Modes Resistively Unstable

- BOUT++ calculations show that C-Mod is *ideal* MHD **stable** for typical EDA H-Modes (1110201023)
- However, such modes become *linearly unstable* when the Pedestal Resistivity is included ($S < 10^9$)

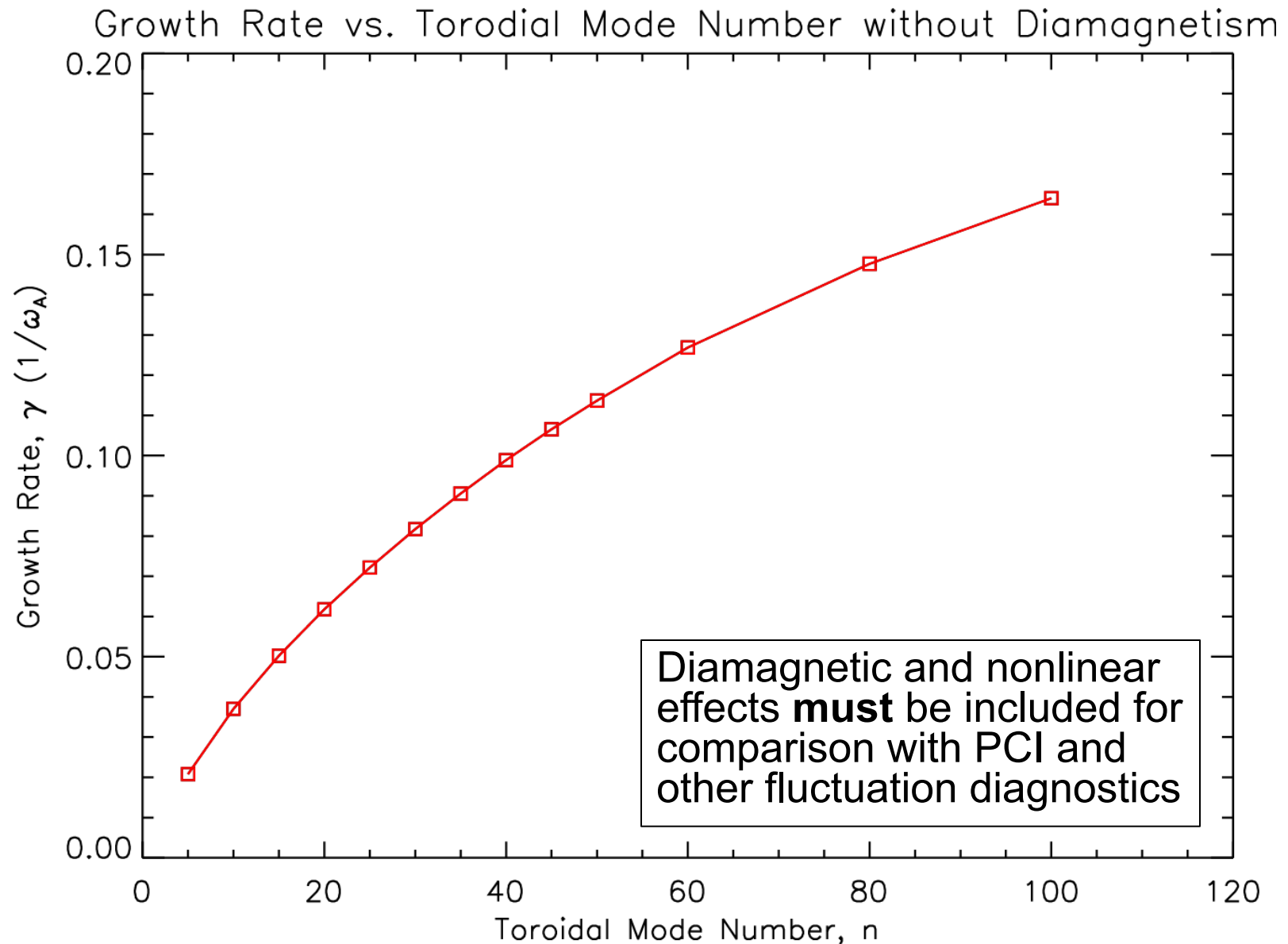


The simulated structure of C-Mod's $n = 15$ resistive ballooning mode

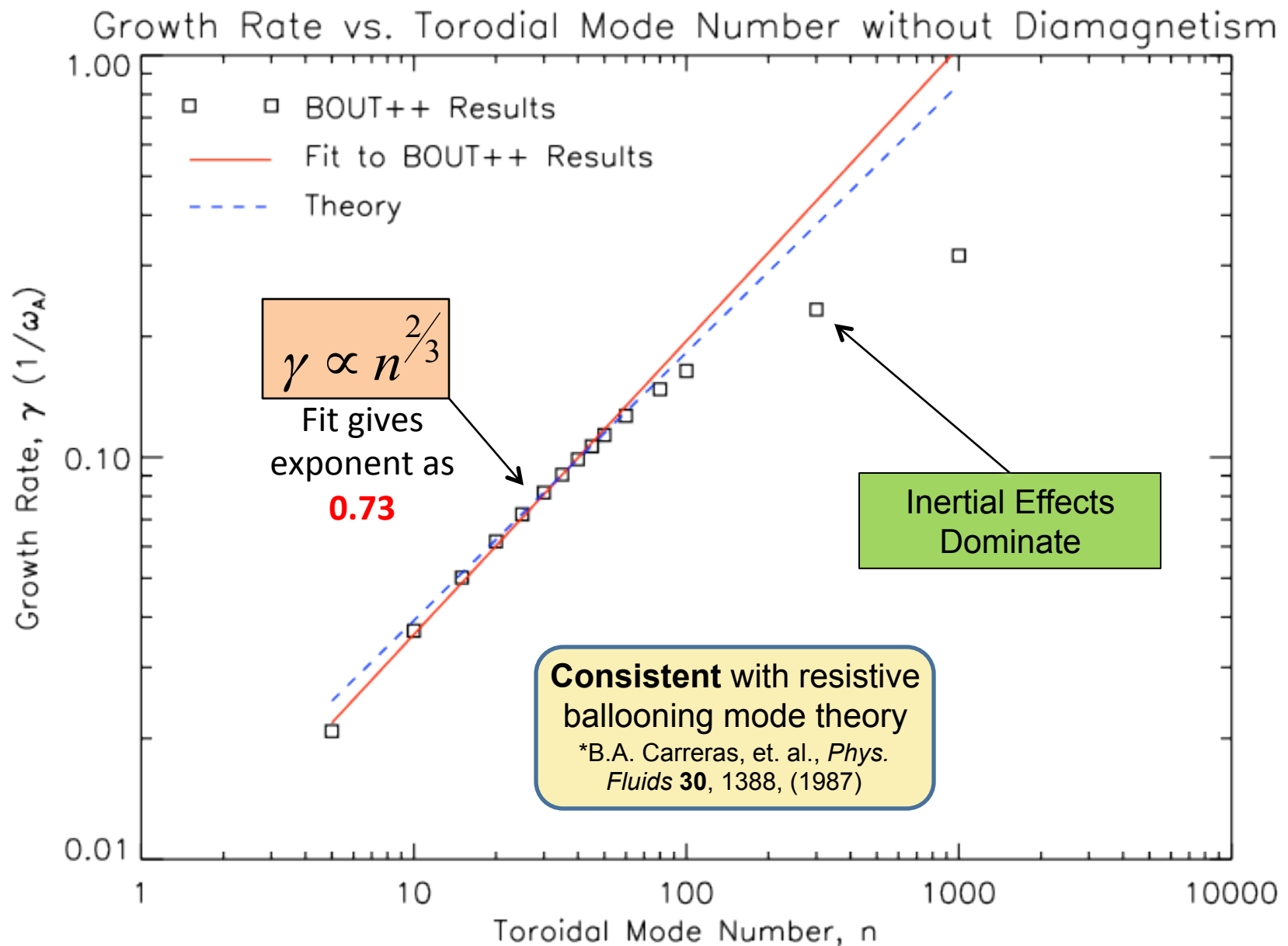
BOUT++ Computed Growth Rates are Consistent with Resistive-Ballooning Mode Theory



BOUT++ has computed the Linear Mode Spectrum for C-Mod's EDA H-Mode



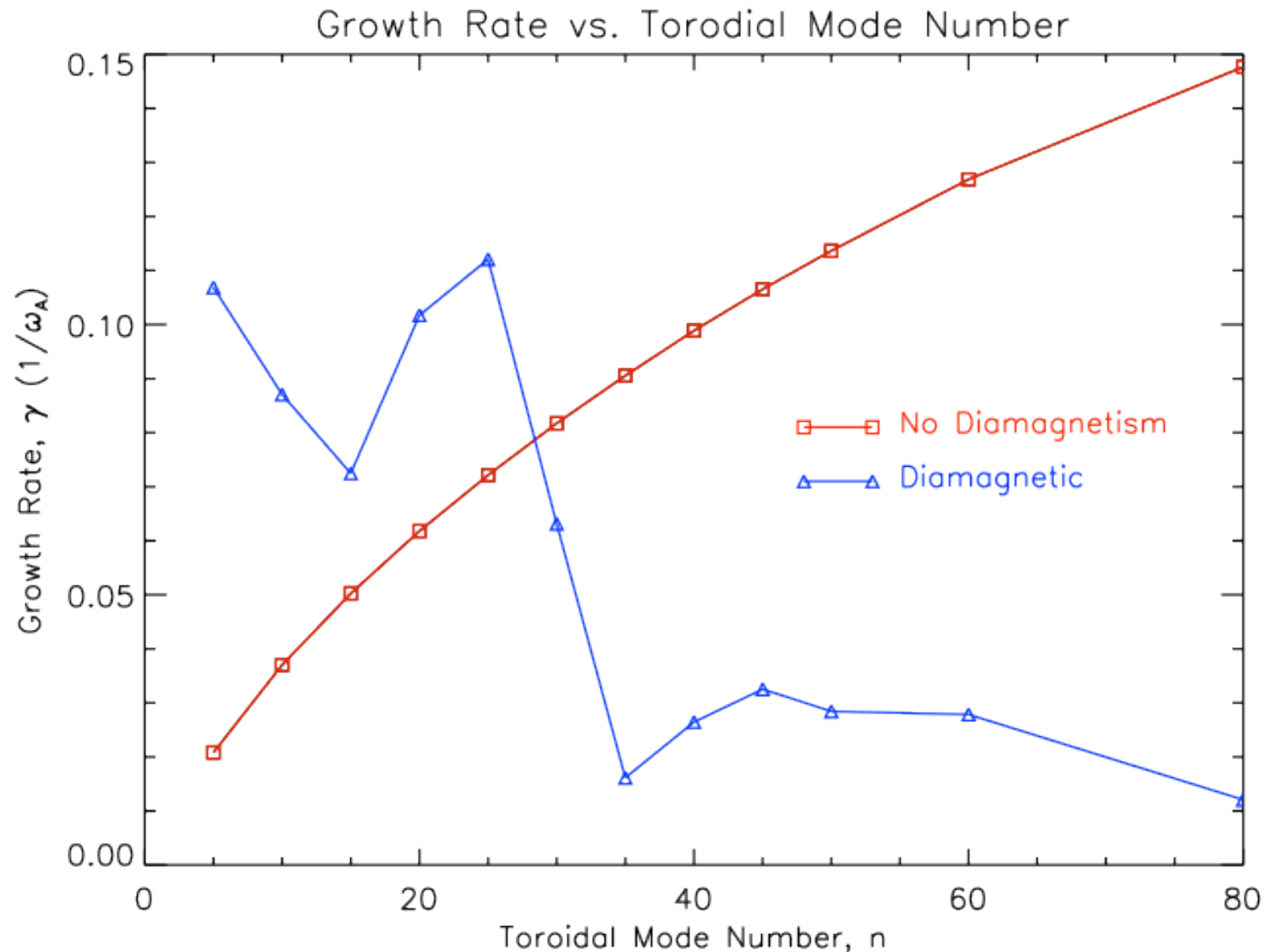
BOUT++ Linear Mode Spectrum Consistent with Resistive-Ballooning Mode Theory



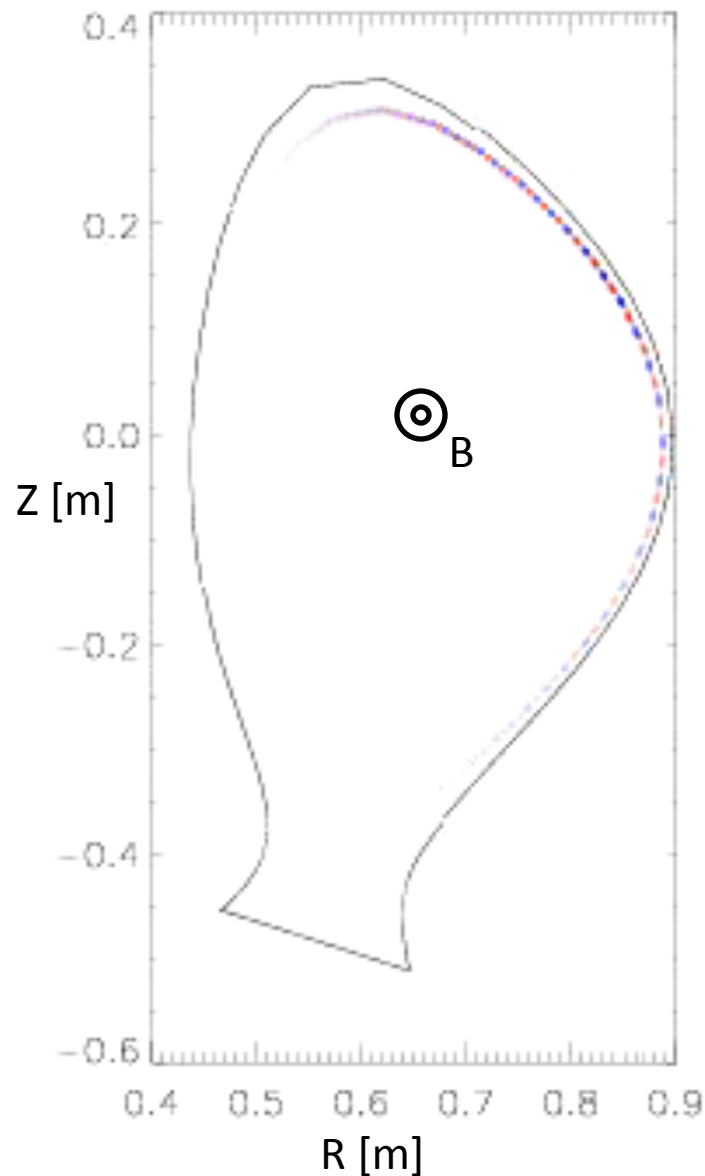
Diamagnetic Effects and an Equilibrium Radial Electric Field were Added into the Model

Reduced MHD Equations	Vorticity	$\frac{\partial \varpi}{\partial t} + v_E \cdot \nabla \varpi = B_0^2 \nabla_{\parallel} \left(\frac{j_{\parallel}}{B_0} \right) + 2b_0 \times \kappa \cdot \nabla p,$	Non-ideal physics ✓ Using resistive MHD term, resistivity can be renormalized as Lundquist Number $S = \mu_0 R v_A / \eta$ ✓ Using hyper-resistivity η_H $S_H = \mu_0 R^3 v_A / \eta_H = S / \alpha_H$ ✓ After gyroviscous cancellation, the diamagnetic drift modifies the vorticity and additional nonlinear terms ✓ Using force balance and assuming no net rotation, $E_{r0} = (1/N_i Z_i e) \nabla_{\perp} P_{i0}$
	Pressure	$\frac{\partial P}{\partial t} + v_E \cdot \nabla P = 0,$	
	Ohm's	$\frac{\partial A_{\parallel}}{\partial t} = -\nabla_{\parallel} (\phi + \Phi_0) + \frac{\eta}{\mu_0} \nabla_{\perp}^2 A_{\parallel} - \cancel{\frac{\eta_H}{\mu_0} \nabla_{\perp}^4 A_{\parallel}},$	
Definitions	$\varpi = \frac{n_0 M_i}{B_0} \left(\nabla_{\perp}^2 \phi + \frac{1}{n_0 Z_i e} \nabla_{\perp}^2 p_i \right), \quad P = P_0 + p$		
	$j_{\parallel} = J_{\parallel 0} - \frac{1}{\mu_0} \nabla_{\perp}^2 A_{\parallel}, \quad v_E = \frac{1}{B_0} b_0 \times \nabla (\phi + \Phi_0)$		

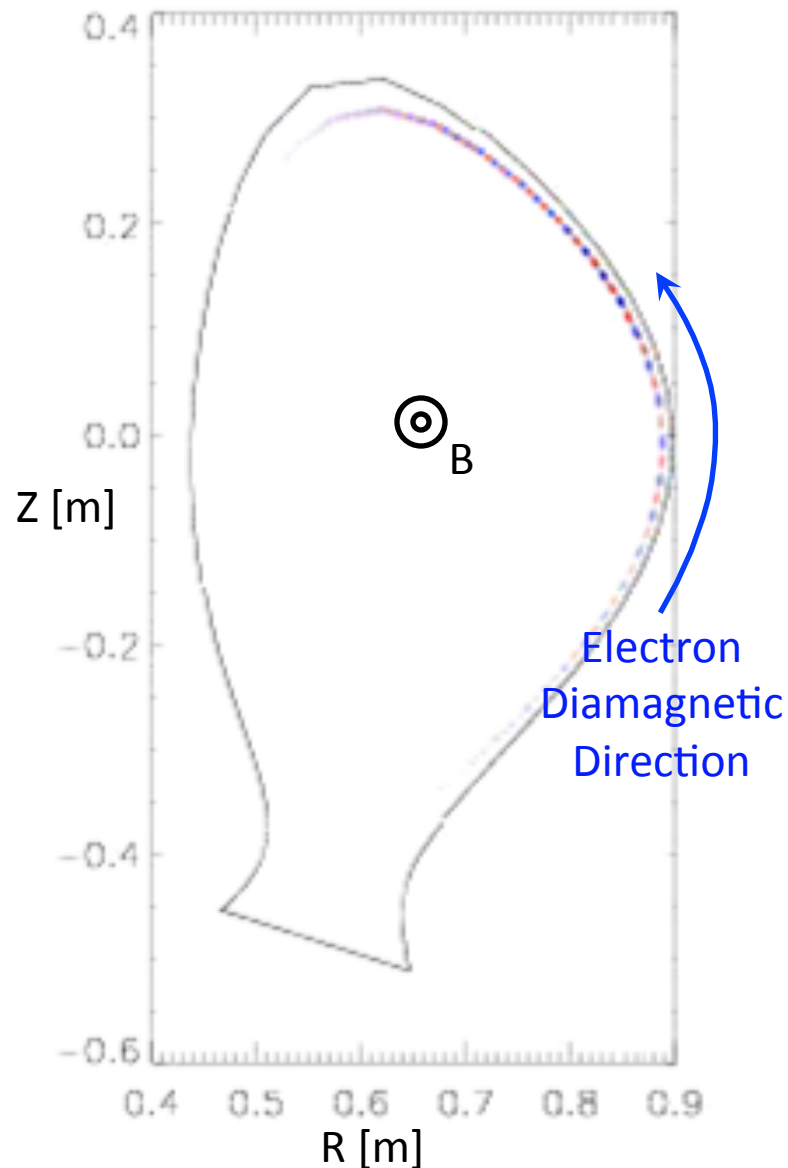
BOUT++ calculations show that Diamagnetic Effects Damp Higher Mode Numbers



The BOUT++ Computed Mode is Found to Propagate in the Electron Diamagnetic Direction

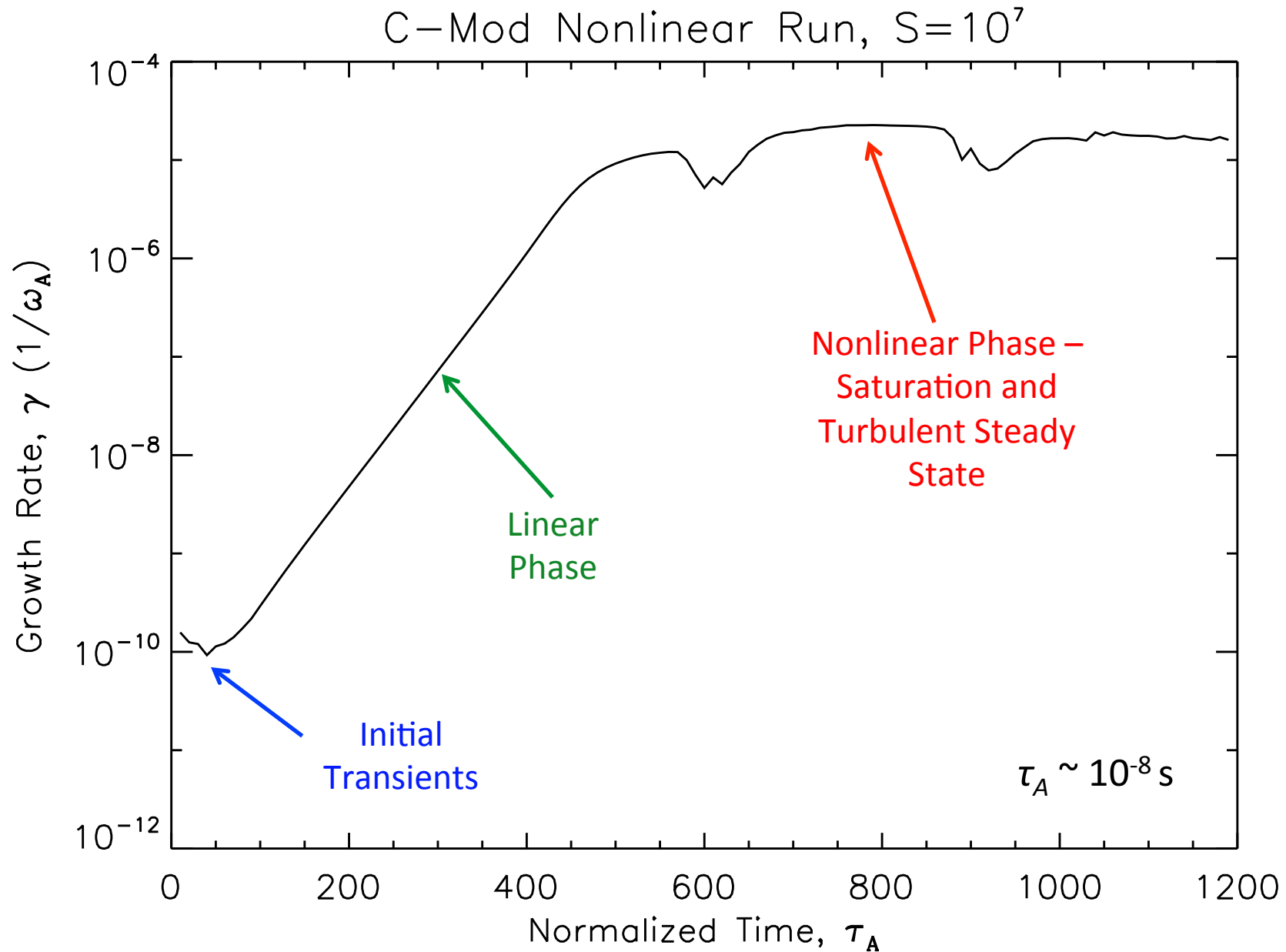


The BOUT++ Computed Mode is Found to Propagate in the Electron Diamagnetic Direction



- Propagation direction agrees with experiment!!!
 - Experimental determination from Correlation between scanning Langmuir probes
- Must run nonlinear simulations to reach saturation and steady-state turbulence
 - Frequency, intensity, and localization of mode can then be compared to measurements
 - PCI, Reflectometry, ECE, etc.
 - Can the QC Mode be excited and controlled by an external antenna?

Preliminary Nonlinear Simulations have begun – Mode Saturation and Turbulent Steady-State have been Observed



Conclusions and Future Work

- BOUT++ results **agree** with theory and show that C-Mod's EDA H-mode is **resistively unstable**
- Turbulent **steady-state** during nonlinear simulations has been achieved
- Incorporating **flow** into nonlinear BOUT++ simulations will allow for comparison with fluctuation diagnostics
 - The physical origins and effects of the EDA QC Mode and the I-mode Weakly Coherent Mode will be investigated
 - This effort will further the understanding of edge turbulence and its influence on tokamak energy confinement